# Perspectives in Low-Energy Nuclear Physics

Robert Roth



TECHNISCHE UNIVERSITÄT DARMSTADT Dear Robert,

David and I would like to ask you if you are willing to give a colloquium at TRIUMF on February 19th 2016.

This is the week before the Nuclear Theory workshop, during which we will have another colloquium by Thomas Papenbrock on "Recent Advances in Nuclear Theory".

In the light of this, we thought that having a preparatory talk where the big questions in nuclear physics are introduced in a very simple manner would be useful...

We feel this will help to communicate the importance of nuclear physics and nuclear theory to the general laboratory...

Please let us know if you are interested in this.

Thanks a lot, Sonia

# The Big Picture

### Low-Energy Nuclear Physics



...a multi-faceted lab connecting fundamental physics with daily life

#### **Diversity**

- structure emerging from microscopic dynamics
- thousands of isotopes, many very special
- zoo of phenomena: shell structure, clustering, deformation, stability, excitations, resonances, transitions, decays, reactions,...

#### Complexity

- multi-physics: strong, weak and el.mag. interactions
- multi-scale: from keV to GeV
- multi-process: structure, decays and reactions
- multi-observable: energies, el.mag.,...

#### Impact

- key input for understanding the evolution of the universe
- tiny things have huge implications: deuteron, Hoyle state,...
- foundation for atomic and molecular physics, chemistry, solid state,...
- direct applications

### Low-Energy Nuclear Theory



...try a non-relativistic quantum description with nucleons as relevant degrees of freedom

### The Problem



## Strategy I: Phenomenology

$$H |\Psi\rangle = E |\Psi\rangle$$

- pick a simple approximation scheme for the Schrödinger equation tailored for the observables of interest
- construct a computationally convenient model for the nuclear interaction compatible with this approximation
- fit the parameters of the interaction model to observables within the chosen approximation scheme

#### **Skyrme-Hartree-Fock**

approximate many-body state by single Slater determinant and parameterize interaction with contact terms

#### Valence-Space Shell Model

freeze core nucleons and treat a few valence nucleons in small valence space with fitted interaction matrix elements

## Strategy II: Ab Initio

$$H |\Psi\rangle = E |\Psi\rangle$$

- construct realistic nuclear interactions based on as much QCD input as possible and fit to few-nucleon properties in exact calculations
- solve the Schrödinger equation using systematic truncations with controlled uncertainties

interactions are independent of many-body framework

different many-body methods can employ the same interaction different interactions can be tested in one many-body approach

### Interactions



~ 1.6fm

 $\rho_0^{-1/3} = 1.8 \text{fm}$ 

nuclear interaction is not fundamental

- residual force analogous to van der Waals interaction between neutral atoms
- based on QCD and induced via polarization of quark and gluon distributions of nucleons
- encapsulates all the complications of the QCD dynamics and the structure of nucleons
- acts only if the nucleons overlap, i.e. at short ranges
- irreducible three-nucleon interactions are important



- first attempts towards construction of nuclear interactions directly from lattice QCD simulations
- compute relative two-nucleon wave function on the lattice
- Invert Schrödinger equation to extract effective two-nucleon potential
- only schematic results so far (unphysical masses and mass dependence, model dependence,...)
- alternatives: phase-shifts or lowenergy constants from lattice QCD

### Tomorrow... from Lattice QCD

Beane et al., PRD87, 034506 (2013), PRC88, 024003 (2013)



# Today... from Chiral EFT

Weinberg, van Kolck, Machleidt, Entem, Meißner, Epelbaum, Krebs, Bernard,...

- low-energy effective field theory for relevant degrees of freedom (π,N) based on symmetries of QCD
- explicit long-range pion dynamics
- unresolved short-range physics absorbed in contact terms, low-energy constants fit to experiment
- systematic expansion in a small parameter with power counting enable controlled improvements and error quantification
- hierarchy of consistent NN, 3N, 4N,... interactions
- consistent electromagnetic and weak operators can be constructed in the same framework



## Many Options

#### standard chiral NN+3N

- NN: N3LO, Entem&Machleidt, nonlocal, cutoff 500 MeV
- 3N: N2LO, Navratil, local, cutoff 500 (400) MeV

#### nonlocal LO...N3LO

• NN: LO...N3LO, Epelbaum, nonlocal, cutoff 450...600 MeV

MeV

3N: N2LO, Nogga, nonlocal, cut

#### N2LO-opt

- NN:
- all these interactions are equally valid, use them all to study uncertainties MeV ., cutoff 500 MeV • 3N: N

#### Iocal N1

- NN: N2LO, Gezerlis et al., local, cutoff 1.0...1.2 fm
- 3N: N2LO, Gezerlis et al., local, cutoff 1.0...1.2 fm

#### semilocal LO...N4LO

- NN: LO...N4LO, Epelbaum, semilocal, cutoff 0.8...1.2 fm
- 3N: N2LO...N3LO, LENPIC, semilocal, cutoff 0.8...1.2 fm

first generation, most widely used up to now

also first generation, but scarcely used

improved fitting, also many-body inputs

designed specifically for QMC applications

order-by-order uncertainty estimates

## Many-Body Solution - Eigenvalue Problem -

## Eigenvalue Problem

start from many-body Schrödinger equation

 $\mathsf{H} | \Psi_n \rangle = E_n | \Psi_n \rangle$ 

define an many-body basis and expand eigenstates in this basis

$$|\Psi_n\rangle = \sum_i C_i^{(n)} |\Phi_i\rangle$$

convert Schrödinger equation into a matrix eigenvalue problem

$$\begin{array}{c} & & \\ & & & \\ & & \\ & & \\ & & \\ & & \\ & & & \\ & & \\ & & & \\ & & \\ & & & \\ & & \\ & & & \\ & & \\ & & & \\ & & \\ & & \\ & & &$$

numerically solve large-scale matrix eigenvalue problem for few low-lying eigenvalues

## **Configuration Interaction**

have to introduce truncations of the basis to make the Hamilton matrix finite and numerically tractable

#### • single-particle truncation (full CI):

truncate the underlying single-particle basis, e.g., at a maximum singleparticle energy e<sub>max</sub>, and construct all possible Slater determinants

#### • many-body truncation (NCSM):

truncate the many-body Slater-determinant basis at a maximum number of harmonic-oscillator excitation quanta  $N_{max}$ 

one has to **demonstrate convergence** with respect to the model-space truncation for accurate results model space dimension **grows dramatically** with truncation parameter  $e_{max}$  or  $N_{max}$ and particle number

incomplete convergence as only source of uncertainties in the many-body treatment

### Convergence



## p-Shell Spectroscopy: Sensitivity

Calci, Roth; arXiv:1601.07209



- of course, spectra depend on choice of initial chiral NN+3N interaction, but the dependence is not dramatic
- important contribution to the total theory uncertainty

## p-Shell Spectroscopy: Sensitivity

Calci, Roth; arXiv:1601.07209



- individual states show systematic disagreement with experiment:
  - second 0<sup>+</sup>: Hoyle state, cluster structure not captured in HO basis
  - first 1<sup>+</sup>: systematic problem with N2LO-3N interaction ?

## p-Shell Spectroscopy: Correlations



- electric quadrupole (E2) observables involving 0<sup>+</sup> ground state & first excited 2<sup>+</sup>
- model-space convergence is terrible, E2 operator sensitive to long-range wave functions

## p-Shell Spectroscopy: Correlations



## p-Shell Spectroscopy: Correlations



## p-Shell Spectroscopy: Precision



## Many-Body Solution - Unitary Transformation -

### **Unitary Transformations**

partially diagonalize Hamilton matrix through a unitary transformation and read-off eigenvalues from the diagonal



continuous unitary transformation of many-body Hamiltonian

 $H_{\alpha} = U_{\alpha}^{\dagger} H U_{\alpha}$ 

morphs the initial Hamilton matrix ( $\alpha = 0$ ) to diagonal form ( $\alpha \rightarrow \infty$ )

## Similarity Renormalization Group

Glazek, Wilson, Wegner, Perry, Bogner, Furnstahl, Hergert, Roth,...

continuous unitary transformation to pre-diagonalize the Hamiltonian with respect to a given basis

consistent unitary transformation of Hamiltonian and observables

$$H_{\alpha} = U_{\alpha}^{\dagger} H U_{\alpha} \qquad O_{\alpha} = U_{\alpha}^{\dagger} O U_{\alpha}$$

design generator for desired diagonalization or decoupling pattern

### In-Medium SRG

Tsukiyama, Bogner, Schwenk, Hergert,...



flow equation for Hamiltonian and Wegner-type generator

$$\frac{d}{ds}H(s) = [\eta(s), H(s)] \qquad \eta(s) = [H(s), H^{\text{off-}'} \text{ additional truncation that causes uncertainties}}$$

write everything in normal ordered form with respect discard normaldeterminantal reference state

$$H(s) = E(s) + \sum_{ij} f_j^i(s) \tilde{A}_j^i + \frac{1}{4} \sum_{ijkl} \Gamma_{kl}^{ij}(s) \tilde{A}_{kl}^{ij} + \frac{1}{36} \sum_{ijklmn} W_{lmn}^{ijk}(s) \tilde{A}_{lmn}^{ljk}$$

## Flowing Energy









## Merging CI and IM-SRG

combine CI/NCSM with IM-SRG to get the best aspects of both methods

solve NCSM problem in small N<sub>max</sub>
extract reference state

- solve MR-IM-SRG flow equations
- decoupling of particle-hole excitations in many-body space
  - solve CI problem with IM-SRG evolved Hamiltonian
  - extract ground and excitation energies and other observables

Gebrerufael et al.; in prep.



Instead of using the zero-body piece of the normal-ordered Hamiltonian, we can use the complete flowing Hamiltonian in a CI calculation

Gebrerufael et al.; in prep.



- Instead of using the zero-body piece of the normal-ordered Hamiltonian, we can use the complete flowing Hamiltonian in a CI calculation
- decoupling though the IM-SRG transformation causes ridiculously fast convergence of CI calculation

Gebrerufael et al.; in prep.



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#### CI-IM-SRG: Excited States

Gebrerufael et al.; in prep.



from the same CI calculation we get the excited states

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- from the same CI calculation we get the excited states
- excitation energies only show subtle changes with IM-SRG flow... ...but there are notable exceptions... Hoyle state?

### CI-IM-SRG: Spectrum

Gebrerufael et al.; in prep.



promising starting point for ab initio studies of arbitrary open-shell nuclei

## Back to Phenomenology

Tsukiyama, et al., PRC 85, 061304(R) (2012); Bogner el al.; PRL 113, 142501 (2014)...

- the decoupling concept can also be used to connect ab initio methods to traditional phenomenological approaches
- valence-space shell model: construct a valence-space interaction by decoupling core and excluded space from valence space though IM-SRG or Lee-Suzuki transformation
- unification of nuclear structure approaches and their interpretation



## New Era of Nuclear Structure Physics

nuclear structure physics has evolved rapidly over the past few years

- many new ideas have emerged for nuclear interactions and many-body approaches... and more is coming
- moving towards a comprehensive theory of nuclear structure and reactions with quantified theory uncertainties
- it's fun to be part of it...

# Epilogue

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